Research Article

# Global regular solutions for the multi-dimensional Kuramoto-Sivashinsky equation posed on smooth domains

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#### **Abstract**

Initial-boundary value problems for the n-dimensional Kuramoto-Sivashinsky equation posed on smooth bounded domains in  $\mathbb{R}^n$  are considered, where n is a natural number from the interval [2,7]. The existence and uniqueness of global regular solutions as well as their exponential decay are established.

Keywords: global solutions; Kuramoto-Sivashinsky equation; decay in bounded domains.

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### 1. Introduction

This work concerns the existence, uniqueness, regularity, and exponential decay rates of solutions to initial-boundary value problems for the n-dimensional Kuramoto-Sivashinsky (KS) equation

$$\phi_t + \Delta^2 \phi + \Delta \phi + \frac{1}{2} |\nabla \phi|^2 = 0. \tag{1}$$

Here n is a natural number from the interval [2,7],  $\Delta$  and  $\nabla$  are the Laplacian and the gradient in  $\mathbb{R}^n$ . In [10], Kuramoto studied the turbulent phase waves and Sivashinsky in [17] obtained an asymptotic equation which simulated the evolution of a disturbed plane flame front (see also [7]). Mathematical results on initial and initial-boundary value problems for the one-dimensional KS equation (1) are presented in [3,5,6,12,14,15,19,20]. The initial-value problem for the multi-dimensional KS type equations (1) was studied in [4,5]. Two-dimensional periodic problems for the KS equation and its modifications posed on rectangles were examined in [2,13,14,16,19], where some results on the existence of weak solutions and nonlinear stability were established. In [11], initial-boundary value problems for the 3D Kuramoto-Sivashinsky equation were studied; the existence, uniqueness, and exponential decay of global regular solutions were proved. For n dimensions,  $x = (x_1, \dots, x_n)$ , n = 2, 3, 4, 5, 6, 7, Equation (1) can be rewritten in the form of the following system:

$$(u_j)_t + \Delta^2 u_j + \Delta u_j + \frac{1}{2} \sum_{i=1}^n (u_i)_{x_j}^2 = 0, \ j = 1, \dots, n,$$
(2)

$$(u_i)_{x_j} = (u_j)_{x_i}, \ j \neq i, \ i, j = 1, \dots, n.$$
 (3)

where  $u_j = (\phi)_{x_j}, \ j = 1, \dots, n$ . Let  $\Omega_n = \prod_{i=1}^n (0, L_i)$  be the minimal nD parallelepiped containing a given smooth domain  $\bar{D}_n$ . The first essential problem that arises while one studies either (1) or (2)–(3), is concerned about the destabilizing effects of  $\Delta u_j$ ; they may be damped by dissipative terms  $\Delta^2 u_j$  provided  $D_n$  has some specific properties. In order to understand this, we use Steklov's inequalities to estimate

$$|a||u_j||^2 \le ||\nabla u_j||^2$$
,  $|a||\nabla u_j||^2 \le ||\Delta u_j||^2$ ;  $a = \sum_{i=1}^n \frac{\pi^2}{L_i^2}$ ,  $j = 1, \dots, n$ .

A simple analysis shows that if

$$1 - \frac{1}{a} > 0,$$
 (4)

then  $\Delta^2 u_j$  damp  $\Delta u_j$ . Naturally, here appear admissible domains where (4) is fulfilled; these are the so-called "thin domains", where some  $L_i$  are sufficiently small while others  $L_j$  may be large  $i, j = 1, \ldots, 7; i \neq j$ .

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The second essential problem is the presence of semi-linear terms in (2) which are interconnected. This does not allow to obtain the first estimate independent of  $u_j$  and leads to a connection between  $L_i$  and  $u_j(0)$ , i, j = 1, ..., 7.

The aim of this paper is to study n-dimensional initial-boundary value problems for (2)–(3) posed on smooth domains, where the existence and uniqueness of global regular solutions as well as their exponential decay of the  $H^2(D_n)$ -norm are established. A "smoothing effect" for solutions with respect to initial data is also obtained. Although, the cases n=2,3 are not new, we included them for the sake of generality.

The remaining part of this paper is organized as follows. Section 2 gives notations and auxiliary facts. In Section 3, formulation of an initial-boundary value problem in a smooth bounded domain  $D_n$  is given. The existence and uniqueness of global regular solutions, exponential decay of the  $H^2(D_n)$ -norm, and a "smoothing effect" are established also in Section 3. Section 4 consists of conclusion.

## 2. Notations and auxiliary facts

Let  $D_n$  be a sufficiently smooth domain in  $\mathbb{R}^n$ , where  $n \in [2,7]$  is a fixed natural number, satisfying the Cone condition (see [1]) and  $x = (x_1, \dots, x_n) \in D_n$ . We use the standard notations of Sobolev spaces  $W^{k,p}$ ,  $L^p$  and  $H^k$  for functions and the following notations for the norms [1] for scalar functions f(x,t):

$$\begin{split} &\|f\|^2 = \int_{D_n} |f|^2 dx, \\ &\|f\|_{L^p(D_n)}^p = \int_{D_n} |f|^p dx, \\ &\|f\|_{W^{k,p}(D_n)}^p = \sum_{0|\leq \alpha|\leq k} \|D^\alpha f\|_{L^p(D_n)}^p, \\ &\|f\|_{H^k(D_n)} = \|f\|_{W^{k,2}(D_n)}. \end{split}$$

When p = 2,  $W^{k,p}(D_n) = H^k(D_n)$  is a Hilbert space with the scalar product

$$((u,v))_{H^k(D_n)} = \sum_{|j| \le k} (D^j u, D^j v), \|u\|_{L^{\infty}(D_n)} = ess \ sup_{D_n} |u(x)|.$$

We use the notation  $H_0^k(D_n)$  to represent the closure of  $C_0^{\infty}(D_n)$ , the set of all  $C^{\infty}$  functions with compact support in  $D_n$ , with respect to the norm of  $H^k(D_n)$ .

**Lemma 2.1** (Steklov's inequality [18]). Let  $v \in H_0^1(0, L)$ . Then

$$\frac{\pi^2}{L^2} \|v\|^2(t) \le \|v_x\|^2(t).$$

**Lemma 2.2** (Differential form of the Gronwall Inequality). Let  $I = [t_0, t_1]$ . Suppose that functions  $a, b : I \to \mathbb{R}$  are integrable and a function a(t) may be of any sign. Let  $u : I \to \mathbb{R}$  be a differentiable function satisfying

$$u_t(t) \le a(t)u(t) + b(t), \text{ for } t \in I \text{ and } u(t_0) = u_0,$$
 (5)

then

$$u(t) \le u_0 e^{\int_{t_0}^t a(t) dt} + \int_{t_0}^t e^{\int_{t_0}^s a(r) dr} b(s) ds.$$

*Proof.* Multiply (5) by the integrating factor  $e^{\int_{t_0}^s a(r) dr}$  and integrate from  $t_0$  to t.

The next Lemmas will be used in estimates.

**Lemma 2.3** (See Theorem 9.1 in [8]). Let n be a natural number from the interval [2,7]. If  $D_n$  is a sufficiently smooth bounded domain in  $\mathbb{R}^n$  satisfying the cone condition and  $v \in H^4(D_n) \cap H^1_0(D_n)$ , then

$$\sup_{D_n} |v(x)| \le C_n ||v||_{H^4(D_n)}.$$

The constant  $C_n$  depends on n and  $D_n$ .

**Lemma 2.4.** Let f(t) be a continuous positive function and f'(t) be a measurable integrable function such that

$$f'(t) + (\alpha - kf^n(t))f(t) \le 0, \ t > 0, \ n \in \mathbb{N},$$
(6)

$$\alpha - kf^{n}(0) > 0, \ k > 0, \tag{7}$$

then

$$f(t) < f(0) \tag{8}$$

for all t > 0.

## 3. KS equation posed on smooth domains

Let  $\Omega_n$  be the minimal nD-parallelepiped containing a given bounded smooth domain  $\bar{D_n} \in \mathbb{R}^n, \ n = 1, \dots, 7$ :

$$\Omega_n = \{x \in \mathbb{R}^n ; x_i \in (0, L_i)\}, \ u_i = (\phi)_{x_i}, \ i = 1, \dots, n.$$

Fix a natural number  $n \in [2,7]$  and consider in  $Q_n = D_n \times (0,t)$  the following initial-boundary value problem:

$$(u_j)_t + \Delta^2 u_j + \Delta u_j + \frac{1}{2} \sum_{i=1}^n (u_i^2)_{x_j} = 0, \ j = 1, \dots, n,$$
(9)

$$(u_i)_{x_i} = (u_j)_{x_i}, \ j \neq i, \ i, j = 1, \dots, n;$$
 (10)

$$u_i|_{\partial D_n} = \Delta u_i|_{\partial D_n} = 0, \ t > 0, \tag{11}$$

$$u_j(x,0) = u_{j0}(x), \ j = 1, \dots, n, \ x \in D_n.$$
 (12)

**Lemma 3.1.** If  $f \in H^4(D_n) \cap H^1_0(D_n)$  and  $\Delta f|_{\partial D_n} = 0$ , then

$$a||f||^2 < ||\nabla f||^2$$

$$a^2 \|f\|^2 < \|\Delta f\|^2$$
.

$$a\|\nabla f\|^2 < \|\Delta f\|^2,$$

$$a^2 \|\Delta f\|^2 \le \|\Delta^2 f\|^2$$

$$\|\Delta \nabla f\|^2 \le \|\Delta^2 f\| \|\Delta f\| \le \frac{1}{a} \|\Delta^2 f\|^2.$$

where

$$a = \sum_{i=1}^{n} \frac{\pi^2}{L_i^2},$$

and

$$||f||^2 = \int_{D_n} f^2(x) dx.$$

Proof. We have

$$\|\nabla f\|^2 = \sum_{i=1}^n \|f_{x_i}\|^2.$$

Define

$$\tilde{f}(x,t) = \begin{cases} f(x,t) & \text{if } x \in D_n; \\ 0 & \text{if } x \in \Omega_n/D_n. \end{cases}$$
(13)

Making use of Steklov's inequalities for  $\tilde{f}(x,t)$  and taking into account that  $\|\nabla f\| = \|\nabla \tilde{f}\|$ , we get

$$\|\nabla f\|^2 \ge a\|f\|^2$$
, where  $a = \sum_{i=1}^n \frac{\pi^2}{L_i^2}$ .

On the other hand,

$$a||f||^2 \le ||\nabla f||^2 = -\int_D f\Delta f dx \le ||\Delta f|| ||f||.$$

This implies

$$a||f|| \le ||\Delta f||$$
 and  $a^2||f||^2 \le ||\Delta f||^2$ .

Consequently,  $a\|\nabla f\|^2 \leq \|\Delta f\|^2$ . Similarly,

$$\|\Delta f\|^2 = \int_{D_n} f\Delta^2 f dx \le \|\Delta^2 f\| \|f\| \le \frac{1}{a} \|\Delta^2 f\| \|\Delta f\|.$$

Hence,  $a\|\Delta f\| \leq \|\Delta^2 f\|$ . Moreover,

$$\|\Delta \nabla f\|^2 = -\int_{D_n} \Delta^2 f \Delta f dx \le \|\Delta^2 f\| \|\Delta f\| \le \frac{1}{a} \|\Delta^2 f\|^2.$$

**Remark 3.1.** Assertions of Lemma 3.1 are true if the function f is replaced respectively by  $u_j$ , j = 1, ..., n.

**Lemma 3.2.** In conditions of Lemma 3.1, the following inequalities hold

$$||f||^2(t)_{H^2(D_n)} \le 3||\Delta f||^2(t),\tag{14}$$

$$||f||^2(t)_{H^4(D_n)} \le 5||\Delta^2 f||^2(t),\tag{15}$$

$$\sup_{D_n} |f(x)| \le C_s ||\Delta^2 f||, \text{ where } C_s = 5C_n.$$
 (16)

*Proof.* To prove (15), we make use of Lemma 3.1 and find

$$||f||_{H^4(D_n)}^2 = ||f||^2 + ||\nabla f||^2 + ||\Delta f||^2 + ||\nabla \Delta f||^2 + ||\Delta^2 f||^2$$

$$\leq \left(\frac{1}{a^4} + \frac{1}{a^3} + \frac{1}{a^2} + \frac{1}{a} + 1\right) ||\Delta^2 f||^2.$$

Since a > 1, then (15) follows. Similarly, (14) can be proved. Moreover, taking into account Lemma 2.3, we get (16).

**Theorem 3.1** (Special basis). Let  $n \in \{2, 3, ..., 7\}$  and  $D_n \in \mathbb{R}^n$  be a bounded smooth domain satisfying the Cone condition. Let  $\Omega_n$  be a minimal nD-parallelepiped containing  $\bar{D_n}$  and

$$\theta = 1 - \frac{1}{a} = 1 - \frac{1}{\sum_{i=1}^{n} \frac{\pi^2}{L_i^2}} > 0.$$
(17)

Given

$$u_{j0}(D_n) \in H^2(D_n) \cap H_0^1(D_n), \ j = 1, \dots, n$$

such that

$$\theta - \frac{2C_s^2 7^3}{a\theta} \left( \sum_{j=1}^n \|\Delta u_j\|^2(0) \right) > 0, \tag{18}$$

then there exists a unique global regular solution to (9)–(12):

$$u_j \in L^{\infty}(\mathbb{R}^+; H^2(D_n)) \cap L^2(\mathbb{R}^+; H^4(D_n) \cap H^1_0(D_n));$$

$$u_{it} \in L^2(\mathbb{R}^+; L^2(D_n)), j = 1, \dots, n.$$

Moreover,

$$\sum_{j=1}^{n} \|\Delta u_j\|^2(t) \le \left(\sum_{j=1}^{n} \|\Delta u_{j0}\|^2\right) \exp\{-a^2 t\theta/2\}$$
(19)

and

$$\sum_{i=1}^{n} \|\Delta u_i\|^2(t) + \int_0^t \sum_{i=1}^{n} \|\Delta^2 u_i\|^2(\tau) d\tau \le C \sum_{i=1}^{n} \|\Delta u_{i0}\|^2, \ t > 0.$$

**Remark 3.2.** In Theorem 3.1, there are two types of restrictions: the first one is pure geometrical,

$$1 - \frac{1}{a} > 0$$

which is needed to eliminate destabilizing effects of the terms  $\Delta u_j$  in (9):

$$||\Delta u_j||^2 - ||\nabla u_j||^2.$$

It is clear that

$$\lim_{L_i \to 0} a = \sum_{i=1}^n \frac{\pi^2}{L_i^2} = +\infty,$$

hence to achieve (17), it is possible to decrease  $L_i$ ,  $i=1,\ldots,n$  allowing other  $L_j$ ,  $j\neq i$  to grow. The situation with condition (18) is more complicated: if initial data are not small, then it is possible either to decrease  $L_i$ ,  $i=1,\ldots,n$ , to fulfill this condition or for fixed  $L_i$ ,  $i=1,\ldots,n$  to decrease initial data  $||u_{i0}||$ .

*Proof of Theorem 3.1.* It is possible to construct Galerkin's approximations to (9)–(12) by the following way. Let  $w_j(x)$  be eigenfunctions of the problem:

$$\Delta^2 w_i - \lambda_i w_i = 0$$
 in  $D_n$ ;  $w_i|_{\partial D_n} = \Delta w_i|_{\partial D_n} = 0, j = 1, 2, \dots$ 

Define

$$u_j^N(x,t) = \sum_{k=1}^N g_k^j(t) w_j(x).$$

Unknown functions  $g_i^j(t)$  satisfy the following initial problems:

$$\left(\frac{d}{dt}u_j^N, w_j\right)(t) + \left(\Delta^2 u_j^N, w_j\right)(t) + \left(\Delta u_j^N, w_j\right)(t) + \frac{1}{2} \left(\sum_{i=1}^n (u_i^N)_{x_j}^2, w_j\right)(t) = 0,$$

$$g_k^j(0) = g_{0k}^j, \quad j = 1, \dots, n, \quad k = 1, 2, \dots$$

The estimates that follow may be established on Galerkin's approximations (see [5,6]), but it is more explicitly to prove them on smooth solutions of (9)-(12).

**Estimate I:**  $u \in L^{\infty}(\mathbb{R}^+; H^2(D_n) \cap H^1_0(D_n)) \cap L^2(\mathbb{R}^+; H^4(D_n) \cap H^1_0(D_n)).$ 

For any natural number  $n \in [2, 7]$ , multiply (9) by  $2\Delta^2 u_i$  to obtain

$$\frac{d}{dt}\|\Delta u_j\|^2(t) + 2\|\Delta^2 u_j\|^2(t) + 2\|\Delta^2 u\|(t)\|\Delta u_j\|(t) + 2\sum_{i=1}^n (u_i(u_i)_{x_j}, \Delta^2 u_j)(t) = 0.$$
(20)

Making use of (15) and Lemmas 2.3, 3.1, 3.2, we write

$$\frac{d}{dt} \|\Delta u_j\|^2(t) + 2\theta \|\Delta^2 u_j\|^2(t) \le 2 \left[ \sum_{i=1}^n \sup_{D_n} |u_i(x,t)| \|\nabla u_i\|(t) \right] \|\Delta^2 u_j\|(t)$$

$$\le 2 \left[ C_s \sum_{i=1}^n \|\Delta^2 u_i\|(t) \|\nabla u_i\|(t) \right] \|\Delta^2 u_j\|(t); j = 1, \dots, n. \tag{21}$$

Summing over j = 1, ..., n and making use of Lemma 3.1, we rewrite (20) in the form:

$$\frac{d}{dt} \sum_{j=1}^{n} \|\Delta u_{j}\|^{2}(t) + 2\theta \sum_{j=1}^{n} \|\Delta^{2} u_{j}\|(t) \leq 2C_{s} n \left( \sum_{j=1}^{n} \|\nabla u_{j}\|(t) \right) \left[ \sum_{j=1}^{n} \|\Delta^{2} u_{j}\|^{2}(t) \right] \\
\leq \left[ \frac{\theta}{2} + \frac{2C_{s}^{2} n^{2}}{\theta} \left( \sum_{j=1}^{n} \|\nabla u_{j}\|(t) \right)^{2} \right] \sum_{j=1}^{n} \|\Delta^{2} u_{j}\|^{2}(t) \\
\leq \left[ \frac{\theta}{2} + \frac{2C_{s}^{2} n^{3}}{\theta} \left( \sum_{j=1}^{n} \|\nabla u_{j}\|^{2}(t) \right) \right] \sum_{j=1}^{n} \|\Delta^{2} u_{j}\|^{2}(t)$$

$$\leq \left[ \frac{\theta}{2} + \frac{2C_s^2 n^3}{a\theta} \left( \sum_{j=1}^n \|\Delta u_j\|^2(t) \right) \right] \sum_{j=1}^n \|\Delta^2 u_j\|^2(t).$$

Taking this into account, we transform (20) in the form

$$\frac{d}{dt} \sum_{j=1}^{n} \|\Delta u_j\|^2(t) + \frac{\theta}{2} \sum_{j=1}^{n} \|\Delta^2 u_j\|^2(t) + \left[\theta - \frac{2C_s^2 n^3}{a\theta} \left(\sum_{j=1}^{n} \|\Delta u_j\|^2(t)\right)\right] \sum_{j=1}^{n} \|\Delta^2 u_j\|^2(t) \le 0.$$
(22)

Condition (18) and Lemma 2.4 guarantee that

$$\theta - \frac{2C_s^2 n^3}{a\theta} \left( \sum_{j=1}^n \|\Delta u_j\|^2(t) \right) > 0, \ t > 0.$$

Hence, (21) can be rewritten as

$$\frac{d}{dt} \sum_{j=1}^{n} \|\Delta u_j\|^2(t) + \frac{a^2 \theta}{2} \sum_{j=1}^{n} \|\Delta u_j\|^2(t) \le 0.$$
 (23)

Integrating, we get

$$\sum_{i=1}^{n} \|\Delta u_j\|^2(t) \le \sum_{j=1}^{n} \|\Delta u_{j0}\|^2 \exp\{-a^2\theta t/2\}.$$
(24)

Consequently, we find

$$\sum_{i=1}^{n} \|\Delta u_i\|^2(t) + \int_0^t \sum_{i=1}^{n} \|\Delta^2 u_i\|^2(\tau) d\tau \le C \sum_{i=1}^{n} \|\Delta u_{i0}\|^2.$$
(25)

Finally, directly from (9), we obtain

$$(u_j)_t \in L^2(\mathbb{R}^+; L^2(D_n)), \ j = 1, \dots, n.$$

Since these inclusions, estimates (24), (25) and Lemma 3.2 do not depend on N, the same estimates are valid also for  $u_j^N(x,t)$ . Hence, it is possible to pass to the limit as  $N \to +\infty$  in

$$u_{j}^{N}(x,t) = \sum_{k=1}^{N} g_{k}^{j}(t)w_{j}(x)$$

and to prove the existence part of Theorem 3.1.

**Lemma 3.3.** A regular solution of Theorem 3.1 is uniquely defined.

*Proof.* Let  $u_j$  and  $v_j$ ,  $j=1,\ldots,n$ , be two distinct solutions to (9)–(12). Denoting  $w_j=u_j-v_j$ , we come to the following system:

$$(w_j)_t + \Delta^2 w_j + \Delta w_j + \frac{1}{2} \sum_{i=1}^n \left( u_i^2 - v_i^2 \right)_{x_j} = 0,$$
(26)

$$(w_j)_{x_i} = (w_i)_{x_i}, \ i \neq j,$$
 (27)

$$w_i|_{\partial D_n} = \Delta w_i|_{\partial D_n} = 0, \ t > 0; \tag{28}$$

$$w_i(x,0) = 0 \text{ in } D, \quad j = 1, \dots, n.$$
 (29)

Multiply (26) by  $2w_i$ , we obtain

$$\frac{d}{dt}\|w_j\|^2(t) + 2\|\Delta w_j\|^2(t) - 2\|\nabla w_j\|^2(t) - \sum_{i=1}^n \left(\{u_i + v_i\}w_i, (w_j)_{x_j}\right)(t) = 0, \ j = 1, \dots, n.$$
(30)

Making use of Lemmas 2.3 and 3.1, 3.2, we estimate

$$I = \sum_{i=1}^{n} \left( \{u_i + v_i\} w_i, (w_j)_{x_j} \right)$$

$$\leq \frac{\epsilon}{2} \|\nabla w_j\|^2 + \frac{1}{2\epsilon} \left( \sum_{i=1}^{n} \|\{u_i + v_i\} w_i\| \right)^2$$

$$\leq \frac{\epsilon}{2a} \|\Delta w_j\|^2 + \frac{2}{\epsilon} \sum_{i=1}^{n} \sup_{D_n} \{u_i^2(x, t)_i + v_i^2(x, t)\} \|w_i^2\|(t)$$

$$\leq \frac{\epsilon}{2a} \|\Delta w_j\|^2 + \frac{2nC_s^2}{\epsilon} \left[ \sum_{i=1}^{n} \{\|\Delta^2 u_i\|^2 + \|\Delta^2 v_i\|^2 \} \right] \sum_{j=1}^{n} \|w_j\|^2.$$

Here,  $\epsilon$  is an arbitrary positive number. Substituting *I* into (30), we get

$$\frac{d}{dt}\|w_j\|^2(t) + (2 - \frac{2}{a} - \frac{\epsilon}{2a})\|\Delta w_j\|^2(t) \le \frac{2nC_s^2}{\epsilon} \left[ \sum_{i=1}^n \{\|\Delta^2 u_i\|^2 + \|\Delta^2 v_i\|^2\} \right] \sum_{j=1}^n \|w_j\|^2.$$
(31)

Taking  $\epsilon = \frac{\theta}{2}$  and summing up over  $j = 1, \dots, n$ , we transform (31) as follows:

$$\frac{d}{dt} \sum_{j=1}^{n} \|w_j\|^2(t) \le C \left[ \sum_{i=1}^{n} \{ \|\Delta^2 u_i\|^2 + \|\Delta^2 v_i\|^2 + \|u_i\|^2(t) + \|v_i\|^2(t) \} \right] \sum_{j=1}^{n} \|w_j\|^2, \ i = 1, \dots, n.$$

By (25) and Lemma 3.1, we have

$$\|\Delta^2 u_i\|^2(t), \|\Delta^2 v_i\|^2(t) \in L^1(\mathbb{R}^+)$$

and

$$||u_i||(t), ||v_i||^2(t) \in L^1(\mathbb{R}^+), i = 1, \dots, n,$$

thence by Lemma 2.2, it holds that

$$||w_j||^2(t) \equiv 0$$
  $j = 1, ..., n$ , for all  $t > 0$ .

Therefore,

$$u_i(x,t) \equiv v_i(x,t); j = 1,\ldots,n.$$

This completes the proofs of Lemma 3.3 and Theorem 3.1.

## 4. Conclusion

This paper is concerned with the formulation and solvability of initial-boundary value problems for the n-dimensional Kuramoto-Sivashinsky system (9)–(10) posed on smooth bounded domains, where  $n \in \{2, 3, ..., 7\}$ . Theorem 3.1 contains results on existence and uniqueness of global regular solutions as well as exponential decay of the  $H^2(D_n)$ -norm, where  $D_n$  is a smooth bounded domain in  $\mathbb{R}^n$ . We define a set of admissible domains, where destabilizing effects of terms  $\Delta u_j$  are damped by dissipativity of  $\Delta^2 u_j$  due to condition (17). This set contains "thin domains", see [9, 13, 16], where some dimensions of  $D_n$  are small while others may be large. Since initial-boundary value problems studied in this paper do not admit a priori estimate independent of  $t, u_j$ , in order to prove the existence of global regular solutions, we put conditions (18) connecting geometrical properties of  $D_n$  with initial data  $u_{j0}$ . Moreover, Theorem 3.1 provides "smoothing effect": initial data  $u_{j0} \in H^2(D_n) \cap H^1_0(D_n)$  imply that  $u_j \in L^2(\mathbb{R}; H^4(D_n))$ .

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## References

- [1] R. A. Adams, Sobolev Spaces, Academic Press, New York, 1975.
- [2] S. Benachour, I. Kukavica, W. Rusin, M. Ziane, Anisotropic estimates for the two dimensional Kuramoto-Sivashinsky equation, *J. Dynam. Differ. Equ.* **26** (2014) 461–476.
- [3] H. A. Biagioni, J. L. Bona, R. J. Iorio, M. Scialom, On the Korteweg-de Vries-Kuramoto-Sivashinsky equation, Adv. Differ. Equ. 1 (1996) 1-20.
- [4] H. A. Biagioni, T. Gramchev, Multidimensional Kuramoto-Sivashinsky type equations: Singular initial data and analytic regularity, *Mat. Contemp.* 15 (1998) 21–42.
- [5] G. Boling, The existence and nonexistence of a global solution for the initial value problem of generalized Kuramoto-Sivashinsky type equations, J. Math. Res. Expo. 11 (1991) 57–69.
- [6] A. T. Cousin, N. A. Larkin, Kuramoto-Sivashinsky equation in domains with moving boundaries, Port. Math. 59 (2002) 335–349.
- [7] M. C. Cross, Pattern formation outside of equilibrium, Rev. Modern Phys. 65 (1993) 851-1086.
- [8] A. Friedman, Partial Differential Equations, Dover, New York, 1997.
- [9] D. Iftimie, G. Raugel, Some results on the Navier-Stokes equations in thin 3D domains, J. Differ. Equ. 169 (2001) 281-331.
- [10] Y. Kuramoto, T. Tsuzuki, On the formation of dissipative structures in reaction-diffusion systems, Prog. Theor. Phys. 54 (1975) 687-699.
- [11] N. A. Larkin, Regularity and decay of solutions for the 3D Kuramoto-Sivashinsky equation posed on smooth domains and parallelepipeds, *Electron. J. Math.* 3 (2022) 1–15.
- [12] N. A. Larkin, Korteweg-de Vries and Kuramoto-Sivashinsky equations in bounded domains, J. Math. Anal. Appl. 297 (2004) 169-185.
- [13] L. Molinet, Local dissipativity in L<sup>2</sup> for the Kuramoto-Sivashinsky equation in spatial dimension 2, J. Dynam. Differ. Equ. 12 (2000) 533–556.
- [14] B. Nicolaenko, B. Scheurer, R. Temam, Some global dynamical properties of the Kuramoto-Sivashinsky equations: nonlinear stability and attractors, Phys. D 16 (1985) 155–183.
- [15] F. Otto, Optimal bounds on the Kuramoto-Sivashinsky equation, J. Funct. Anal. 257 (2009) 2188–2245.
- [16] G. R. Sell, M. Taboada, Local dissipativity and attractors for the Kuramoto-Sivashinsky equation in thin 2D domains, Nonlinear Anal. 18 (1992) 671–687.
- [17] G. I. Sivashinsky, Nonlinear analysis of hydrodynamic instability in laminar flames-1: Derivation of basic equations, *Acta Astronauica* 4 (1977) 1177–1206.
- [18] A. V. Steklov, The problem of cooling of an heterogeneous rigid rod, Commun. Kharkov Math. Soc. 5 (1896) 136–181.
- [19] R. Temam, Infinite-Dimensional Dynamical Systems in Mechanics and Physics, Springer, New York, 1988.
- [20] J. Li, B. Y. Zhang, Z. Zhang, A nonhomogeneous boundary value problem for the Kuramoto-Sivashinsky equation in a quarter plane, *Math. Methods Appl. Sci.* 40 (2017) 5619–5641.